# ON THE NONHOLONOMIC COMPONENTS OF THE CHRISTOFFEL SYMBOLS IN $V_n$ (I)

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V. 공간에서 Christoffel symbol의 Non-holonomic components에 관하여 (I)

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### Summary

In this paper we study the inverse of the results obfained in our previous paper; Hyun, J.O & Kim, H.G 1980(on the christo ffel symbols of the non-holonomic frames in  $V_{\pi}$ ) in order to reconstruct and to investigate the useful relationships between holononomic and nonholonomic components of the christoffel symbols.

## 1. INTRODUCTION.

Let  $h_{\lambda\mu}$  be the fundamental metric tensor, whose determinant

(1.1)  $h \stackrel{\text{def}}{=} \det(h_{\lambda\mu}) \neq 0$ and let  $e^{\nu}$   $(i=1,2,\ldots,n)$  be a set of *n* linearly independent vectors in *n*-dimensional Riemannian space  $V_n$  refered to a real coordinate system  $x^{\nu}$ .

Then there is a unique tensor  $h^{i\nu} = h^{\nu_i}$  defined by

$$(1.2) h_{\lambda\mu} h^{\lambda\nu} \stackrel{\text{def}}{=} \delta^{\nu}_{\mu}$$

and a unique reciprocal set of *n* linearly independent covariant vectors  $e_i^i$  (*i*=1,2,..., *n*,), satisfying

$$(1.3)^* \qquad e^{\nu} \stackrel{i}{e_{\lambda}} = \delta^{\nu}_{\lambda} , \quad e^{\lambda} \stackrel{i}{e_{\lambda}} = \delta^{i}_{j}.$$

Within the vectors  $e^{r}$  and  $e_{i}$  a nonholonomic frame of  $V_{\pi}$  defined in the following way;

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If  $T_{a}^{****}$  are holonomic components of a tensor, then its nonholonomic components are defined by

(1.4) a 
$$T_{j}^{i} \cdots \stackrel{\text{def}}{=} T_{j}^{i} \cdots \stackrel{e_{j}}{=} e_{j}^{i} \cdots$$
  
From (1.3) and (1.4) a  
(1.4) b  $T_{j}^{i} \cdots \stackrel{\text{def}}{=} T_{j}^{i} \cdots \stackrel{e_{j}}{=} e_{k}^{i} \cdots$ .

In this paper, for our further discussion, results obtained in our previous paper Chung, K.T & Hyun, J.O 1976 and Hyun, J.O & Kim, H.G 1980 will be introduced without proof.

### 2. PRELIMINARY RESULTS.

Theorem (2.1). We have

<sup>(\*)</sup> Throughout the present paper, all indices take the values 1,2,..., *n* and follow: the summation convention. Greek indices are used for the holonomic components of a tensor, while Roman indices are used for the nonholonomic components of a tensor.

#### 2 논 문 집

(2.1) a  $e^* = e_{\lambda} h_{ij} h^{\lambda *}, \quad e^j_{\lambda} = e^* h^{ij} h_{\lambda *}$ (2.1) b  $h_{ij} = \delta_{ij}, \quad h^{ij} = \delta^{ij}, \quad e^* = e^*, \quad e^j_{\lambda} = e_{\lambda}$ 

Consider a symmetric covariant tensor a whose determinant  $a \stackrel{\text{def}}{=} \det (a_{2\mu}) \neq 0$ . It is well-known that the quantities defined by

# $a^{\lambda^{\nu}} \stackrel{\text{def}}{=} \frac{\text{cofector of } a_{\lambda_{\nu}} \text{ in } a}{a}$

is a symmetric contravariant tensor satisfying

$$(2.2) a_{\lambda\mu} a^{\lambda\nu} = \delta^{\nu}_{\mu}$$

**Theorem** (2.3). The holonomic and nonholonomic components of the christoffel symbols satisfy

(2.3) a 
$$[jk, m]_a = [\lambda \mu, \omega]_a \stackrel{a}{j} \stackrel{a}{k} \stackrel{a}{m} \stackrel{a}{m} + a_{\lambda\mu}(\partial_r e^{\lambda}) \stackrel{e^{\lambda}}{k} \stackrel{e^{\mu}}{m} \stackrel{e^{\mu}}{k}$$
  
(2.3) b  $\left\{ \begin{array}{c} i\\ jk \end{array} \right\}_a = - \left\{ \begin{array}{c} \nu\\ \lambda \mu \end{array} \right\}_a \stackrel{e^{\lambda}}{e} \stackrel{e^{\lambda}}{p} \stackrel{e^{\lambda}}{k} \stackrel{e^{\lambda}}{m} \stackrel{e^{\lambda}}{m} \stackrel{e^{\lambda}}{k} \stackrel{e^{\lambda}}{m} \stackrel{e^{\lambda}$ 

Here,  $[jk, m]_e$  and  $\left\{ \begin{array}{c} i\\ jk \end{array} \right\}_e$  are the christoffel symbols of the first and second kind, respectively defined by  $a_{ie}$ .

**Theorem** (2.3). The nonholonomic components of the christoffel symbols of the second kind may be expressed as

(2.4). 
$$\left\{ \begin{array}{c} i\\ jk \end{array} \right\}_{a} = -e^{a} e^{a} \quad a \quad e_{a}$$

Where  $\nabla_{\mu}$  is the symbol of the covariant derivative with respect to  $\begin{cases} \nu \\ \lambda_{\mu} \end{cases}_{\mu} \end{cases}_{\mu}$ .

In this section, we consider the inverse of the obtained previous results, reconstruct and in ve stigate the relationships between the holonomic and nonhalonomic components of the christoffel symbols.

# 3. HOLONOMIC AND NONHOLON-OMIC COMPONENTS OF CHRI-STOFFEL SYMBOLS IN V<sub>s</sub>.

Let  $a_{2p}$  and  $a_{ij}$  are holonomic and nonholonomic components of the tensor and take a coordinate system  $y^i$  for which we have at a point p of  $V_n$ 

(3.1) a 
$$\frac{\partial y_i}{\partial x^4} = e_{i}, \quad \frac{\partial x^*}{\partial y^i} = e_{i}^*.$$

We have

**Theorem** (3.1). The holonomic components of the christoffel symbols, as follows;

(3.2) a 
$$[\lambda \mu, \omega]_{a} = [jk, m]_{a} \overset{j \quad k}{e_{l}} \overset{m}{e_{l}} \overset$$

**Proof.** From (1.4) b,

$$(3,3) a_{\lambda\mu} = a_{jk} e_{\lambda} e_{\mu}.$$

Differentiating with respec to the coordinate system  $x^{\bullet}$  of  $V_{\pi}$ . We have

$$(3.4) \qquad \partial_{\bullet}(a_{\lambda \mu}) = \partial_{m}(a_{jk}) \stackrel{j \quad k \quad m}{e_{\lambda} \quad e_{\lambda} \quad e_{\mu}} \stackrel{j \quad k \quad m}{e_{\lambda} \quad e_{\lambda}} \stackrel{j \quad k \quad m}{e_{\lambda} \quad e_{\lambda} \quad e_{\lambda}}$$

The following equation (3.5) a is obtain from (3.4) by interchanging  $\omega$  and  $\mu$ , *m* and *k* throughout, (3.5) by interchaging  $\omega$  and  $\lambda$ , *m* and *j*;

(3.5) a 
$$\partial_{\mu}(a_{\lambda \nu}) = \partial_{k}(a_{jm}) \stackrel{j}{e_{\lambda}} \stackrel{m}{e_{\nu}} \stackrel{k}{e_{\nu}} \stackrel{j}{e_{\nu}} \stackrel{m}{e_{\nu}} + a_{jm}(\partial_{\mu} \stackrel{j}{e_{\nu}}) \stackrel{m}{e_{\nu}} \stackrel{j}{e_{\nu}} + a_{jm}e_{\lambda} (\partial_{\mu} \stackrel{m}{e_{\nu}})$$

- 252 -

(3.5) b 
$$\partial_{\lambda}(a_{\mu\nu}) = \partial_{j} (a_{km})e_{\mu} e_{\nu} e_{\lambda}$$
  
  $+ a_{km}(\partial_{\lambda} e_{\mu}) e_{\nu} e_{\nu}$   
  $+ a_{km}e_{\mu} (\partial_{\lambda} e_{\nu})$ 

The sum of (3.5)a and (3.5)b substract (3.4)and divide by 2, and by means of (3.3), we have the first relation (3.2)a as in following ways;

(3.6) 
$$[\lambda \mu, \omega]_{a} = [jk, m]_{a} e_{\lambda} e_{\mu} e_{a}$$
$$+ a_{jk} (\partial_{\mu} e_{\lambda}) e_{a}$$
$$= [jk, m]_{a} e_{\lambda} e_{\mu} e_{a}$$
$$+ a_{jk} (\partial_{\mu} e_{\lambda}) e_{\mu} e_{a}.$$

. Multiplying both sides of (2.3) a by  $e_a e_r e_r$ , according to (1.4) a and

 $(3.7) \qquad \begin{array}{c} e^{\theta} = e^{\lambda} & \delta^{\theta}.\\ k & k & \lambda \end{array}$ 

We have the same results as (3.6).

The second relation (3.2) b may be obtain by multiplying  $e^{a}_{i} e_{\beta} e_{\tau}$  to both sides of (2.3) b and using (1.3) and (2.2), (3.7)

$$\begin{cases} i \\ jk \end{cases}_{a} e^{\alpha} e^{\beta} e^{\beta} e^{\gamma}_{T} = \begin{cases} \nu \\ \lambda \mu \end{cases}_{a} e^{\nu} e^{\lambda} e^{\beta} e^{\alpha} e^{\beta} e^{\alpha}_{T} \\ i \\ e^{\nu} e^{\alpha} e^{\beta} e^{\beta} e^{\gamma}_{T} (\partial_{\mu} e^{\nu}) \\ e^{\beta} e^{\beta} e^{\gamma}_{T} (\partial_{\mu} e^{\alpha}) e^{\beta} e^{\beta} \\ e^{\beta} e^{\gamma}_{T} e^{\beta} e^{\beta} e^{\gamma}_{T} (\partial_{\mu} e^{\alpha}) e^{\alpha}_{T} \\ e^{\beta} e^{\gamma}_{T} e^{\beta} e^{\beta} e^{\gamma}_{T} e^{\beta} e^{\beta} e^{\alpha}_{T} e^{\beta} e^{\beta} e^{\beta} e^{\alpha}_{T} e^{\beta} e^{\beta} e^{\beta} e^{\alpha}_{T} e^{\beta} e^{\beta} e^{\beta} e^{\gamma}_{T} e^{\beta} e^$$

**Theorem** (3.2). The holonomic components of the christoffel symbols of the second kind may be expressed as

(3.8) 
$$\begin{cases} \alpha \\ \beta \gamma \end{cases}_{a} = -e_{\beta}^{j} e_{r}^{k} \left( \nabla_{k} e_{j}^{e} \right)$$
$$= e_{r}^{k} e_{j}^{e} \left( \nabla_{k} e_{\beta}^{i} \right)$$

Where  $\nabla_k$  is the symbol of the covariant derivative with respect to  $\begin{cases} i\\ jk \end{cases}_a$ 

**Proof.** Using (2.1) a and (3.7), We have (3.8) from (3.2) b as in the following way;

$$\begin{cases} \alpha \\ \beta \gamma \end{cases}_{a} = \begin{cases} i \\ jk \end{cases}_{a} e^{\alpha} e^{j} e^{k} e_{j} e_{r} - (\partial_{\tau} e^{\alpha}) e^{j} e_{r} \\ = \begin{cases} i \\ jk \end{cases}_{a} e^{\alpha} e^{j} e_{r} e_{r} (\partial_{m} e^{\alpha}) e^{j} e_{r} e_{r} \\ = -e^{j} e^{k} e_{r} (\partial_{k} e^{\alpha} - \begin{cases} i \\ jk \end{cases}_{a} e^{\alpha} e^{\alpha}$$

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4 논 문 집

〈국문초록〉

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V. 공간에서 Christoffel symbol의 Non-holonomic components에 관하여 (1)

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본 논문은 앞 논문(현진오·김홍기, 1980)에서 얻어진 결과의 역을 증명함으로써 christoffel symbol의 holonomic과 nonholonomic component 사이의 관계를 더욱더 명확히 하고 이에 대한 효율적이고 새로운 표 현방법을 연구했다.

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