A Thesis for the Degree of M.E.

A Note on the Curl of the Nonholonomic Vector in V_n

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감사의 글

이 논문이 완성되기까지 바쁘신 가운데도 많은 지 도를 아끼지 않으신 현진오교수님께 감사 드리며,아 울러 그동안 많은 도움을 주신 수학교육과의 여러 교 수님과 동료들에게 심심한 사의를 표합니다. 그리고, 그 동안 저에게 사랑과 격려를 하여 주신 가족,친지 및 주위의 많은 분들께 또한 감사를 드립니다.



1985년6월 일

강 성 홍

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ABSTRACT (KOREAN)



I. INTRODUCTION

Introducing a set of 4 linearly independent basic null vectors, V.Hlavaty(Geometry of Einstein's Unified Field Theory. P. Noordhoff Ltd. 1957) introduced the concept of the nonholonomic frames and used it successfully as a tool to develop the algebra of the unified field theory in the space-time X_{\pm}

In the present papers, the curl of the vector a_{λ} will be proved in a refined way as application of orthogonal nonholonomic frames.

Let V_n be a n-dimensional Riemannian space referred to a real coordinate system x^{ν} and defined by fundamental metric tensor $h_{\lambda\mu}$, whose determinant is an SOLAR

(1.1) h $\frac{\text{def}}{\text{Det}(h_{\lambda\mu})} \neq 0$.

Then there is the unique tensor $h^{\lambda \nu} = h^{\nu \lambda}$ defined by

$$(1.2) \quad h_{\lambda\mu} h^{\lambda\nu} \frac{\det \delta^{\nu}}{\mu}.$$

Consider a set of n linearly independent vectors e^{ν} (i = 1,2 ..., ., n). There is the unique reciprocal set of n linearly independent covariant vectors e_{λ}^{j} (j = 1, 2, ..., n) satisfying

(1.3)
$$e^{\nu} \dot{e}_{\lambda}^{j} = \delta_{\lambda}^{\nu} **$$

With these vectors e^{ν} and e^{λ}_{λ} , a nonholonomic frame of V_{n} may be constructed in the following way: If T_{λ}^{ν} are holonomic components of a tensor, the nonholonomic components of the holonomic tensor T_{λ}^{ν} are defined by

(1.4)
$$T_{j}^{j}$$
 $\underline{\det f} T_{\lambda}^{j}$ $e_{\nu}^{j} e_{j}^{\lambda}$

and

(1.5) $T^{\nu}_{\lambda} \dots \underline{def}^{*} T^{j}_{j} \dots \underline{e}^{\nu} e^{\lambda}_{j} \dots$



In this section, for our further discussion, results obtained in our previous papers will be introduced without proof.

THEOREM 2.1. The covariant derivative of the holonomic covariant vector is given by

(2.1)
$$\begin{array}{c} \nabla_{\mu}(a_{\lambda}) = \left[\partial_{k}^{*}a_{j} - a_{j}^{*}\right]_{jk}^{k} = \mu^{k}a_{\lambda} \\ = \nabla_{k}^{*}a_{j} - \mu^{k}a_{\lambda} \\ \end{array}$$
 where $\begin{array}{c} \nabla_{\mu}a_{\lambda} = \partial_{\mu}a_{\lambda} - a_{\nu} \\ \nabla_{\lambda\mu} \end{array}$.

^{**.} Throughout the present paper, Greek indices take values 1,2,n unless explicitly stated otherwise and follow the summation convention, while koman indices are used for the nonholonomic components of a tensor and run from 1 to n. Roman indices also follow the summation convention.

THEOREM 2.2. The covariant derivative of the nonholonomic covariant vector is equivalent to

(2.2)
$$V_{k}^{*} a_{j} = \left[\partial_{\mu} a_{\lambda} - a_{\nu} \int_{\lambda \mu}^{\nu} \right] e^{\lambda} e^{\mu} e^{\mu}$$

$$= V_{\mu} a_{\lambda} e^{\lambda} e^{\mu} e^{\mu}.$$

COROLLARY 2.3. We have

(2.3)
$$V_{\mu}a_{\lambda} = \partial_{\mu}a_{\lambda} - a_{j}(V_{\mu}e_{\lambda}).$$

THEOREM 2.4. We have

 $(2.4) \quad \nabla_{\mu} a_{\lambda} = \partial_{k}^{*} a_{j} e_{\lambda}^{j} e_{\mu}^{k} + a_{j}^{*} (\nabla_{\mu} e_{\lambda}^{j}).$

THEOREM 2.5. The covariant derivative of the holonomic covariant tensor $a_{\mu\lambda}$ may be expressed in terms of the nonholo-nomic components;

(2.5)
$$V_{\mu} a_{\nu\lambda} = \left[\partial_{k}^{*}a_{jj} - a_{\ell j}^{*}\int_{ik}^{\ell}a_{i\ell} \int_{kj}^{\ell}b_{j}\right] \stackrel{j}{}_{\nu} \stackrel{k}{}_{\lambda} \stackrel{j}{}_{\mu} \stackrel{k}{}_{\nu} \stackrel{j}{}_{\lambda} \stackrel{k}{}_{\mu}$$

THEOREM 2.6. We have

(2.6)
$$P_{k}^{*} a_{jj} = \left[\partial_{\mu} a_{\nu\lambda} - a_{\omega\lambda}\right]^{-\omega} - a_{\nu\omega}\left[\mu\lambda\right] e^{\nu} e^{\lambda} e^{\mu}$$

 $i j k$

.THE CURL OF THE VECTOR

In this section, we shall reconstruct the curl of a vector and obtain its special properties with holonomic and nonholonomic frames. -3-

THEOREM 3.1. For the curl of the vector a_{λ} , following four expressions are equal to each other.

(3.1) (a)
$$\nabla_{\mu} a_{\lambda} - \nabla_{\lambda} a_{\mu}$$

(b) $\partial_{\mu} a_{\lambda} - \partial_{\lambda} a_{\mu}$
(c) $(\partial_{k} a_{j} - \partial_{j}^{*} a_{k}) e_{\lambda} e_{\mu}$
(d) $(\nabla_{k} a_{j} - \nabla_{j}^{*} a_{k}) e_{\lambda} e_{\mu}$.

PROOF. The equality of (a) and (b) is given by the

result

$$(3.2) \quad \nabla_{\mu} a_{\lambda} = \partial_{\mu} a_{\lambda} - a_{\lambda} - a_{\nu} \int_{\lambda \mu}^{\nu} d\mu$$

The equality of (b) and (c) is given by (2.1). By means of (2.2), (c) is equal to (d).

COROLLARY 3.2. The curl of the nonholonomic vector a j may be expressed in terms of the components holonomic curl.

PROOF. Using (2.2), we have

$$(3.3) \qquad \nabla_{\mathbf{k}}^{*} \mathbf{a}_{\mathbf{j}} - \nabla_{\mathbf{j}}^{*} \mathbf{a}_{\mathbf{k}} = \left[\nabla_{\mu} \mathbf{a}_{\lambda} - \nabla_{\lambda} \mathbf{a}_{\mu} \right] \mathbf{e}_{\mathbf{k}}^{\mu} \mathbf{e}_{\mathbf{j}}^{\lambda}.$$

THEOREM 3.3. If A is the curl of a covariant vector, we have the following equations;

(3.4) (a)
$$\nabla_{\nu} A_{\lambda\mu} + \nabla_{\lambda} A_{\mu\nu} + \nabla_{\mu} A_{\nu\lambda} = 0$$

(b) $\partial_{\nu} A_{\lambda\mu} + \partial_{\lambda} A_{\mu\nu} + \partial_{\mu} A_{\nu\lambda} = 0$
 $-4-$

PROOF . By means of the curl,

$$(3.5) \qquad \nabla_{\nu} A_{\lambda\mu} = \nabla_{\nu} \left[\nabla_{\mu} a_{\lambda} - \nabla_{\lambda} a_{\mu} \right]$$

$$\nabla_{\lambda} A_{\mu\nu} = \nabla_{\lambda} \left[\nabla_{\nu} a_{\mu} - \nabla_{\mu} a_{\nu} \right]$$

$$\nabla_{\mu} A_{\nu\lambda} = \nabla_{\mu} \left[\nabla_{\lambda} a_{\nu} - \nabla_{\nu} a_{\lambda} \right]$$

Using the properties of the covariant derivative, the sum of the left hand side of three equations is identically zero. The covariant derivative of holonomic tensor $A_{\lambda\mu}$ is given by

$$\begin{split} F_{\nu} A_{\lambda\mu} &= \partial_{\nu} A_{\lambda\mu} - A_{\sigma\mu} \int_{\lambda\nu}^{\sigma} - A_{\lambda\sigma} \int_{\nu\mu}^{\sigma} \\ F_{\lambda} A_{\mu\nu} &= \partial_{\lambda} A_{\mu\nu} - A_{\sigma\nu} \int_{\mu\lambda}^{\sigma} - A_{\mu\sigma} \int_{\lambda\nu}^{\sigma} \\ F_{\mu} A_{\nu\lambda} &= \partial_{\mu} A_{\nu\lambda} - A_{\sigma\lambda} \int_{\nu\mu}^{\sigma} - A_{\nu\sigma} \int_{\mu\lambda}^{\sigma} . \end{split}$$

Since $A\nu_{\lambda}$ is skew-symmetric, we have

THEOREM 3.4. Let *A is be a curl of the nonholonomic covariant vector a_{λ} . Then the covariant derivative of nonholonomic tensor *A is may be expressed as following relation;

$$(3.6) \qquad \nabla_{j}^{*} A_{jk} + \nabla_{j}^{*} A_{kj} + \nabla_{k}^{*} A_{jj}$$
$$= \partial_{j}^{*} A_{jk} + \partial_{j}^{*} A_{kj} + \partial_{k}^{*} A_{jj} = 0$$

PROOF. Making use of (2.5), we obtain

(3.7)
$$\nabla_{i}^{*} A_{jk} = \partial_{i}^{*} A_{jk} - A_{\ell k}^{*} \int_{ji}^{\ell} A_{j\ell} \int_{\ell}^{\ell} \int_{\ell}^{\ell} A_{k} \int_{\ell}^{\ell} \int_{\ell}^{\ell} A_{k} \int_{\ell}^{\ell} \int_{\ell}^{\ell} A_{k} \int_{$$

By similar method

$$(3.8) \qquad \nabla_{j}^{*} A_{kj} = \partial_{j}^{*} A_{kj} - A_{\ell i} \int_{kj}^{\ell} - A_{k\ell} \int_{ij}^{\ell} \partial_{k} \partial_{k} \partial_{j} \partial_{k} \partial_{$$

Since *A_{ij} is skew-symmetric, the sum of each side of (3.7)

and (3.8) is
$$\bigwedge_{j=0}^{\infty} \bigwedge_{j=0}^{\infty} \bigwedge_{$$

On the other hand, from (2.2)

$$(3.10) \qquad \nabla_{i}^{*} A_{jk} + \nabla_{j}^{*} A_{ki} + \nabla_{k}^{*} A_{jj}$$
$$= \left(\nabla_{\nu} A_{\lambda\mu} + \nabla_{\lambda} A_{\mu\nu} + \nabla_{\mu} A_{\nu\lambda} \right) \qquad e^{\nu} e^{\lambda} e^{\mu}_{k} .$$

But the first term of right side of (3.10), by means of (3.4), is zero.

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<國文抄錄>

V, 空間에서의 NONHOLONOMIC VECTOR의 CURL에 對하여

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本 論文의 重要한 目的은 RIEMANN 空間 V_n에서 HOLONOMIC 構造를 갖는 VECTOR의 CURL에 對한 몇 가지 性質들을 NONHOLONOMIC 構造를 利用하여 보다 새로운 方法으로 證明하는데 있다.

