ON SOME RELATIONS OF TWO NONHOLONOMIC CHRISTOFFEL SYMBOLS IN Vn

By

Oh, Bong Leem

Major in Mathematics Graduate School of Education Jeju National University

Supervised By

Assistant Prof. Hyun, Jinoh

May, 1983

ON SOME RELATIONS OF TWO NONHOLONOMIC CHRISTOFFEL SYMBOLS IN Vn

이를 教育學碩士學位 論文으로 提出함



濟州大學校教育大學院數學教育專攻

提出者 吳 奉 琳

指導教授 玄 進 五

1983年 5月 日

吳奉琳의 碩士學位 論文을 認准함

濟州大學校教育大學院 主審 副審 副審

1983年 5月 日

감사의 글

이 논문이 완성되기까지 연구에 바쁘신 가운데도 자상한 지도를 하여 주신 현진오 교수님께 무한한 감 사를 드리오며 그동안 많은 도움을 주신 수학교육과의 모든 교수님께 심심한 사의를 표합니다. 그리고 그동안 저에게 사랑과 격려를 아끼지 않으신 주위의 여러분들께 감사 드립니다.

1983년 5월 일

오 봉 립

CONTENTS

ABSTRACT (KOREAN)

1.	INTRODUCTION	1
2.	PRELIMINARY RESULTS	3
3.	MAIN RESULTS 제주대학교 중앙도서관	5
	LITERATURE CITED	7
	ABSTRACT(ENGLISH)	8

국 문 초 록

Vn 공간에서 두개의 NONHOLONOMIC CHRIS -

TOFFEL 기호의 어떤 관계에 대하여

제주대학교 교육대학원

수학교육전공

오 봉 림

본 논문의 목적은 일반적인 대칭공변 TENSOR 아/ 에 의하여 정의된 CHRISTOFFEL기호에 대한 HOLONOMIC 과 NONHOLONOMIC성분들 사이의 관계를 연구하고, 제1과 제2 NONHOLONOMIC CHRISTOFFEL 기호 들에 대한 유용한 표현 형식을 유도하였다.

1. INTRODUCTION

Let V_n be a *n*-dimensional Riemannian space referred to a real coordinate system x^* and defined by a fundamental metric tensor h_{1n} , whose determinant

(1.1) $h \stackrel{\text{def}}{=} \text{Det} ((h_{\lambda \mu})) \neq 0.$

According to (1.1), there is a unique tensor $h^{\lambda\nu} = h^{\nu\lambda}$ defined by

 $(1.2) \qquad h_{\lambda\mu} h^{\lambda\nu} \frac{\mathrm{def}}{\mathrm{def}} \delta^{\nu}_{\mu}.$

Let e^{v} (i = 1, 2, ..., n) be a set of *n* linearly independent vectors. Then there is a unique reciprocal set of *n* linearly independent covariant vector e^{i}_{i} (i = 1, 2, ..., n), satisfying

(1.3) $e_i^{*} e_i^{i} = \delta_{\mu}^{*}$ (*) With these vectors e_i^{*} and e_i^{i} a nonnolonomic frame of V_n may

with these vectors e and e, a nonnoronomic frame of v, may

be constructed in the following way:

If A_{λ}^{sc} :: are holonomic components of a tensor, then its nonholonomic components are defined by

(1.4) $A_j^i \dots = A_{\mu}^{\nu} \dots e_{\nu} e_{\nu}^{\mu}$

From the above definition, we obtain that

$$(1.5) \qquad A_{\lambda}^{i} \dots = A_{j}^{i} \dots e_{\lambda}^{i} e_{\lambda} \dots$$

(*) Throughout the present paper, Greek indices take the values 1,2, ..., n unless explicitly stated otherwise and follow the summation convention, while Roman indices are used for the nonholonomic components of a tensor and run from I to n. Roman indices also follow the summation convention.

With respect to orthogonal nonholonomic frame of V_n constructed by an orthogonal ennuple e^v ($i = 1, 2, \dots, n$) it was shown by Chung. K.T.& Hyun, J.O. 1976 that

$$(1.6) h_{ij} = \delta_{ij} , h^{ij} = \delta^{ij}$$

(1.7)
$$e^{\nu} = e^{\nu}, e_{\lambda} = e_{\lambda},$$

In this paper, studying the relationships between holonomic and nonholonomic components of the Christoffel symbols defined by a general symmetric covariant tensor $a_{\lambda\mu}$, we derive a useful representation of the first and second nonholonomic Christoffel symbols.



2. PRELIMINARY RESULTS

Consider a symmetric covariant tensor a whose determinant $a \underline{\det}$ Det $(a_{\lambda\mu}) \neq 0$. It is well known that the quantities $a^{\lambda\nu}$ defined by

 $a^{\lambda\nu} \underline{\det} \quad \underline{\det} \quad \underline{cofactor \ of \ a_{\lambda\nu} \ in \ a}_{a}$ is symmetric contravariant tensor satisfying

 $(2.1) a_{\lambda\mu} a^{\lambda\nu} = \delta^{\nu}_{\mu}.$

Take a coordinate system y^i for which we have at a point P of V_n

(2.2)
$$\frac{\partial y^{i}}{\partial x^{\lambda}} = \stackrel{i}{e}_{\lambda}, \quad \frac{\partial x^{\nu}}{\partial y^{i}} = \stackrel{i}{e}_{i}^{\nu}$$

If $a_{\lambda\mu}$ and a_{ij} are holonomic and nonholonomic components of the tensor defined above, it follows that

(2.3)
$$a_{jk} = a_{\lambda\mu} e^{\lambda} e^{\mu}$$
. TIONAL UNIVERSITY LIBRARY

In this section, for our further discussions, results obtained in the previous paper will be introduced without proof.

THEOREM 2.1. The derivative of e^{λ} is a negative self-adjoint. That is

(2.4) $\partial_k \left(\begin{array}{c} i \\ e_{\lambda} \end{array} \right) e^{\mu} = -\partial_k \left(\begin{array}{c} e^{\mu} \\ j \end{array} \right) \stackrel{i}{e_{\lambda}}.$

THEOREM 2.2. The derivative of the tensor $a_{\lambda\mu}$ is a negative self-adjoint.

 $(2.5) a^{\lambda\mu} \partial_k (a_{\lambda\mu}) = -a_{\lambda\mu} \partial_k (a^{\lambda\mu}).$

THEOREM 2.3. The holonomic and nonholonomic components of the Christoffel symbols satisfy

-3-

$$(2.6) \qquad \{m, jk\}_{a} = \{\omega, \lambda\mu\}_{a} e^{\lambda} e^{\mu} e^{\omega}, \\ + a_{\lambda\mu} (\partial_{\tau} e^{\lambda}) e^{\tau} e^{\mu}, \\ (2.7) \qquad \{\frac{i}{jk}\}_{a} = \{\lambda\mu\}_{a} e^{i}_{\nu} e^{\lambda} e^{\mu} + e^{i}_{\nu} e^{\mu} (\partial_{\mu} e^{\nu}).$$

.

Here, [m, jk] and $\{\frac{i}{jk}\}$ are the first and second Christoffel symbol of nonholonomic frame, respectively, defined by $a_{\lambda\mu}$.

THEOREM 2.4 The nonholonomic components of the Christoffel symbols of the second kind may be expressed as

$$(2.8) \qquad \{ \begin{array}{l} i\\ j\\ k \end{array} \}_{a} = \begin{array}{l} e\\ e\\ u\\ k \end{array} \stackrel{\mu}{=} \left(\begin{array}{c} \partial\\ \mu\\ k \end{array} \stackrel{\mu}{=} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \stackrel{\mu}{=} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = \begin{array}{c} e\\ i\\ k \end{array} \right) \\ = - \begin{array}{c} e\\ e\\ i \end{array} \stackrel{\mu}{=} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ e\\ i \end{array} \stackrel{\mu}{=} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} \partial\\ \mu\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \bigg) \\ = - \left(\begin{array}{c} e\\ i \end{array} \bigg(\left(\begin{array}{c} e\\ i \end{array} \right) \\ = - \left(\begin{array}{c} e\\ i \end{array} \bigg) \\ = - \left(\begin{array}{c} e\\ i \end{array}$$

Where ∇_{μ} is the symbol of the covariant derivative with respect to $\{ \begin{array}{c} \nu \\ \lambda \mu \end{array} \}_{a}$

THEOREM 2.5 The holonomic components of the Christoffel symbols, as follows

$$(2.9) \qquad (\omega, \lambda \mu)_{a} = (m, jk)_{a} \stackrel{m}{e}_{\omega} \stackrel{j}{e}_{\lambda} \stackrel{k}{e}_{\mu} + a_{jk} (\partial_{\mu} \stackrel{j}{e}_{\lambda}) \stackrel{k}{e}_{\omega}.$$

$$(2.10) \qquad \{\stackrel{\alpha}{\beta}_{r}\}_{a} = \{\stackrel{i}{jk}\}_{a} \stackrel{e^{\alpha}}{e}_{\beta} \stackrel{j}{e}_{r} - (\partial_{r} \stackrel{e^{\alpha}}{e}_{\beta}) \stackrel{j}{e}_{\beta}$$

$$= \{\stackrel{i}{jk}\}_{a} \stackrel{e^{\alpha}}{e}_{\beta} \stackrel{j}{e}_{r} + (\partial_{r} \stackrel{j}{e}_{\beta}) \stackrel{e^{\alpha}}{j}$$

$$= -\stackrel{j}{e}_{s} \stackrel{k}{e}_{r} (\nabla_{k} \stackrel{j}{e}_{\beta})$$

$$= \stackrel{k}{e}_{r} \stackrel{e^{\alpha}}{e}_{\beta} (\nabla_{k} \stackrel{j}{e}_{\beta}).$$

3. MAIN RESULTS.

THEOREM 3.1 Ine nonholonmic components of the Christoffel symbols of the first and second kinds satisfy

(3.1) $\{ \begin{array}{c} i \\ jk \\ a \end{array} = a^{im} (m, jk)_{a}.$

PROOF. From(2.7),(2.9)

$$\begin{cases} p \\ qr \end{cases}_{a} = \{ \frac{\lambda}{\mu\omega} \}_{a} \stackrel{p}{e_{\lambda}} e^{\mu} e^{\omega} + \stackrel{p}{e_{\lambda}} e^{\omega} (\partial_{\omega} e^{\lambda}) \\ = a^{\lambda\sigma} (\sigma, \mu\omega) a \stackrel{p}{e_{\lambda}} e^{\mu} e^{\omega} + \stackrel{p}{e_{\lambda}} e^{\omega} (\partial_{\omega} e^{\lambda}) \\ = a^{\lambda\sigma} ((i, jk) a \stackrel{i j k}{e^{\sigma} e^{\mu} e^{\omega}} + a_{jk} (\partial_{\omega} e^{\mu}) e_{\sigma}) \\ \stackrel{p}{e_{\lambda}} e^{\mu} e^{\omega} + \stackrel{p}{e_{\lambda}} e^{\omega} (\partial_{\omega} e^{\lambda}) \\ e_{\lambda} e^{\mu} e^{\omega} + \stackrel{p}{e_{\lambda}} e^{\omega} (\partial_{\omega} e^{\lambda}) \\ \{ p \\ qr \}_{a} = a \stackrel{i p}{[i, jk]} \stackrel{j k}{\delta} \stackrel{j k}{\delta} + a \stackrel{p k}{a} \stackrel{j k}{a_{jk}} (\partial_{\omega} e_{\mu}) e^{\mu} e^{\omega} \\ + \stackrel{p}{a_{i}} e^{\omega} (\partial_{\omega} e^{\lambda}) \\ = a \stackrel{i p}{[i, qr]}_{a} + (\partial_{\omega} \stackrel{p}{e_{\mu}}) e^{\mu} e^{\omega} - (\partial_{\omega} \stackrel{p}{e_{\lambda}}) e^{\lambda} e^{\lambda} e^{\omega} \\ = a \stackrel{i p}{[i, qr]}_{a} . \end{cases}$$

By the theorem 3.1.

COROLLARY 3.2. We have

$$a_{il}\left\{\frac{i}{jk}\right\}_{a}=\left\{l,jk\right\}_{a}$$

PROOF. Multiplying both sides of (3.1) by a_{il} , using (2.1) as required.

-5-

THEOREM 3.3 The nonholonomic components of the Christoffel symbols of the second kind may be expressed as the symbols of the covariant derivative with respect to the second kind.

PROOF. Using(2.7),(2.10),

$$\{ \begin{array}{l} i\\ jk \}_{a} = \{ \begin{array}{l} \lambda\\ \mu\omega \}_{a} \begin{array}{l} i\\ e_{\lambda} e^{\mu} e^{\omega} + e^{\omega} + e^{\lambda} e^{\omega} (\partial_{\omega} e^{\lambda}) \\ j\\ k\\ \end{array} \right. \\ = \begin{array}{l} m\\ e_{\omega} e^{\lambda} (\mathbf{\nabla}_{m} e^{\mu}) \\ i\\ e_{\lambda} e^{\mu} e^{\omega} + e^{\lambda} e^{\omega} (\partial_{\omega} e^{\lambda}) \\ j\\ k\\ \end{array} \right. \\ = \left(\mathbf{\nabla}_{k} e^{\mu} \right) e^{\mu}_{j} + \begin{array}{l} i\\ e_{\lambda} e^{\omega} (\partial_{\omega} e^{\lambda}) \\ k\\ \end{array} \right)$$

where $abla_{p}$ is the symbol of the covariant derivative with

respect to {ⁱ_{jk}}_a 제주대학교 중앙도서관 JEDU NATIONAL UNIVERSITY LIBRARY

LITERATURE CITED

- C.E.Weatherburn, 1957. An Introduction to Riemannian Geometry and the tensor calculus. Jambridge University Press.
- Chung K.T. & Hyun J.O. 1976. On the nonholonomic frames of
 V_n, Yonsei Nonchong. Vol. 13.
- o. J.O.Hyun & H.G.Kim. 1981. On the Christoffel symbols of the Nonholonomic Frames in V_{n} .
- o. J.O.Hyun & E.S.Bang. 1981. On the Nonholonomic components of the Christoffel symbols in $V_n(I)$
- o. J.O.Hyun & T.C.Kang. A note on the Nonholonomic self- Adjoint in $\ensuremath{V_n}$

.

ABSTRACT

ON SOME RELATIONS OF TWO NONHOLONOMIC CHRISTOFFEL SYMBOLS IN V_{m}

Oh, Bong Leem Department of Mathematics Graduate School of Education Jeju National University

The purpose of the present paper is to study the relationships between holonomic and nonholonomic components of the Christoffel symbols defined by general symmetric covariant tensor $a_{\lambda\mu}$ and we derive a useful representation of the first and second nonholonomic Christoffel symbols.

-8-