On the Christoffel Symbols of the Non-holonomic fromes in Vn

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RIEMANN공간 Vn에서의 비-호로노미 구조에 대한 CHRISTOFFEL SYMBOL에 과학 소고

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Summary

The purpose of the present paper is primarily to study the relationships between holonomic and nonholonomic components of the christoffel symbols defined by a general symmetric covariant tensor $a_{\lambda\mu}$. Secondarily, we derive a useful representation of the christoffel symbols, formed with respect to the metric tensor $h_{\lambda\mu}$ of V_{μ} , in terms of orthogonal ennuple.

I. INTRODUCTION

Let V_n be a *n*-dimensional Riemannian space referred to a real coordinate system x^{ν} and defined by a fundamental metric tensor $h_{\lambda\mu}$, whose determinant

(1.1)
$$h \stackrel{\text{def}}{=} Det(h_{1\mu}) \neq 0.$$

According to (1.1) there is a unique tensor $h^{\lambda^{\nu}} = h^{\nu_{\lambda}}$ defined by

$$(1.2) \qquad h_{\lambda\mu} h^{\lambda\nu} \stackrel{\text{def}}{=} \delta^{\mu}_{\mu}$$

The tensors $h_{2\mu}$ and $h^{2\nu}$ will serve for raising and lowering indices of tensor quantities in V_{π} in the usual manner. Let e^{i} , (i=1,...,n), be a set of *n* linearly independent vectors. Then there is a unique reciprocal set of *n* linearly independent covariant vectors e_{i} (i=1,...,n), satisfying

(1.3)a
$$e^{*} e_{\lambda} = \delta^{*}_{\mu}^{(**)}$$

With these vectors e^{v} and e_{1}^{i} a nonholonomic frame of V_{n} may be constructed as in the follow ing way: If T_{μ}^{*} :::: are holonomic components of a tensor, then its nonholonomic components are defined by

(1.4)a
$$T_{j\cdots}^{i\cdots} \stackrel{\text{def}}{=} T_{j\cdots}^{i\cdots} \stackrel{i}{e}, e^{i}_{j\cdots}$$

An easy inspection of (1,3)a and (1,4)a shows

^(*) Supported by the research fund of the Ministry of Education, R.O.K., 1979.

^(**) Throughout the present paper, all indices take the values 1, 2, ... n and follow the summation convention. Greek indices are used for the holonomic components of a tensor, while Roman indices are used for the nonholonomic components of a tensor.

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that

(1.4)b
$$T_{\lambda...}^{*...} = T_{j...e^{*}}^{*...e^{*}} e_{\lambda}^{*...}$$

With respect to an orthogonal nonholonomic frame of V_n constructed by an orthogonal emmple e_i^* , $(i=1,\ldots,n)$, it was shown in Chung, K.T & Hyun, J. O. 1976. that

(1.5)a
$$h_{ij} = \partial_{ij}, h^{ij} = \partial^{ij};$$

(1.5)b $e^{\nu} = e^{\nu}, e_{1} = e_{1}, i$

The purpose of the present paper is, in the first, to study the relationships between holonomic and nonholonomic components of the Christoffel symbols defined by a general symmetric covariant tensor $\sigma_{2\mu}$. In the second, we derive a useful representation of Christoffel symbols, formed with respect to the metric tensor $h_{2\mu}$ of V_{π} in terms of orthogonal ennuple.

I. NONHOLONOMIC COMPONENTS OF CHRISTOFFEL SYMBOLS IN V.

Consider a symmetric covariant tensor a whose determinant $a \det Det((a_{\lambda_{\mu}}))$ is not zero. It is well known that the quantities $a^{\lambda_{\mu}}$ defined by

$$a^{*} \operatorname{\underline{def}}_{a} \operatorname{\underline{cofactor of } a_{2}, in a}_{a}$$

is a symmetric contravariant tensor satisfying

$$(2.1) a_{\lambda\mu} a^{\lambda\nu} = \delta_{\mu}^{\nu}.$$

Theorem (2.1). The holonomic and nonholonomic components of Christoffel symbols satisfy

(2.2)a
$$[jk, m]_{e} = [\lambda \mu, \omega]_{e} \frac{e^{\lambda}e^{\mu}e^{\mu}}{e^{\lambda}\mu} + a_{2\mu}(\partial_{T}e^{\lambda}) \frac{e^{\tau}e^{\mu}}{e^{\lambda}\mu},$$

(2.2)b
$$\begin{cases} \frac{i}{jk} \\ \frac{i}{j} \\ e^{\lambda}\mu \end{cases} = \begin{cases} \nu \\ \lambda \mu \end{pmatrix} \frac{i}{j} \frac{e^{\lambda}e^{\mu}}{e^{\lambda}\mu} + \frac{e^{\lambda}e^{\mu}}{e^{\lambda}\mu} (\partial_{\mu}e^{\nu}) \end{cases}$$

Here, $[jk, m]_{a}$ and $\begin{cases} j \\ jk \end{cases}_{a}$ are the Christoffel symbols of the first and second kind, respectively, defined by a_{1a} .

Proof. Take a coordinate system y^i for which we have at a point P of V_n

(2.3)
$$\frac{\partial y^i}{\partial x^\lambda} = e_\lambda, \quad \frac{\partial x^\nu}{\partial y^i} = e_\lambda^\nu.$$

If $a_{\lambda\mu}$ and a_{ij} are holonomic and nonholonomic components of the tnesor defined above, it follows from (1.4) a that

$$(2.4) a_{jk} = a_{\lambda\mu} e^{\lambda} e^{\mu}$$

Differentiating with respect to y", we have

(2.5)
$$\partial_m a_{jk} = (\partial_u a_{\lambda\mu}) e^{\lambda} e^{\mu} e^{\nu} + a_{\lambda\mu} (\partial_m e^{\lambda}) e^{\mu} e^{\mu} + a_{\lambda\mu} e^{\lambda} (\partial_m e^{\mu}).$$

The first of the following equations is obtained from (2.5) by interchanging k and m throughout and the dummy indices ω and μ , the second by interchanging j and m and the dummy indices λ and ω :

$$\partial_{\underline{k}} a_{j\underline{m}} = (\partial_{\mu} a_{\lambda \omega}) \underbrace{e^{\lambda} e^{\omega} e^{\mu}}_{j \ \underline{m} \ \underline{k}} + a_{\lambda \omega} (\frac{\partial e^{\lambda}}{\partial t}) \underbrace{e^{\omega}}_{\underline{k} \ \underline{m}} + a_{\lambda \omega} \underbrace{e^{\lambda} (\partial_{\underline{k}} e^{\omega})}_{\underline{k} \ \underline{m}},$$

$$(2.6)$$

$$\partial_{\lambda} a_{\underline{k}\underline{m}} = (\partial^{\lambda} a_{\mu \omega}) \underbrace{e^{\mu} e^{\omega} e^{\lambda}}_{\underline{k} \ \underline{m} \ \underline{m}} + a_{\mu \omega} (\partial_{j} e^{\omega}) \underbrace{e^{\mu}}_{\underline{m} \ \underline{k}} + a_{\mu \omega} (\partial_{j} e^{\mu}).$$

If from the sum of these two equations in (2.6) we subtract (2.5) and divide by 2, we have in consequence of (2.3) the first relation (2.2)a as in the following way:

$$[jk, m]_{e} = [\lambda \mu, \omega]_{e} e^{\lambda e^{\mu}} e^{\mu} + a_{\lambda \mu} e^{\mu} (\partial_{\mu} e^{\lambda})$$
$$= [\lambda \mu, \omega]_{e} e^{\lambda e^{\mu}} e^{\mu} + a_{\lambda \mu} (\partial_{\tau} e^{\lambda}) e^{\tau} e^{\mu}$$

The second relation (2.2) may be obtained by multiplying

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a^{im}=a^{aβ}e_eep

to both sides of (2, 2)a and by making use of (1, 3) and (2, 1), as in the following way:

Theorem (2.2). The nonholonomic components of the Christoffel symbols of the second kind may be expressed as

(2.7)
$$\left\{ \begin{array}{c} i\\ jk \end{array} \right\}_{a} = - e^{v} e^{\mu} \nabla^{i}_{\mu} e_{\nu},$$

where ∇_{μ} is the symbol of the covariant derivative with respect to $\left\{ \begin{array}{c} \nu \\ \lambda \mu \end{array} \right\}_{a}$ **Proof.** Making use of (1.3)a, we may derive (2.7) from (2.2)b as in the following way:

$$\begin{cases} i\\ jk \end{cases}_{e} = \stackrel{i}{e_{v}} e^{\mu} (\partial_{\mu} e^{\nu} + \begin{cases} \nu\\ \lambda \mu \end{cases}_{a} e^{\lambda}) = \stackrel{i}{e_{v}} e^{\mu} \nabla_{\mu} e^{\nu} \\ = - \frac{e^{\nu}}{j} e^{\mu} (\nabla_{\mu} e^{i}_{v}). \end{cases}$$

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〈초 록〉

RIEMANN 공간 Vn에서의 비-호로노미 구조에 대한 Christoffel Symbol에 관한 소고

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본 논문에서는 Riemann 공간 Vn에서의 일반적인 Symmetric Covariant tensor a₂₄에 의하여 정의되어진 Christoffel Symbol의holonomic과 nonholonomic compoment 사이의 관계를 구명하고 이에 대한 효율적이**교** 새로운 표현방법을 연구했다.

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